

ELECTRON ENERGY SPECTRUM AND MAGNETIC INTERACTIONS  
IN HIGH- $T_c$  SUPERCONDUCTORS

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The character of magnetic interactions in La-Sr-Cu-O and Y-Ba-Cu-O systems is of primary importance for analysis of high- $T_c$  superconductivity in these compounds. Neutron diffraction experiments showed the antiferromagnetic ground state for nonsuperconducting  $\text{La}_2\text{CuO}_4$  and  $\text{YBa}_2\text{Cu}_3\text{O}_8$  with the strongest antiferromagnetic superexchange being in the ab plane [1,2]. Nonsuperconducting "1-2-3" system has even two Neel temperatures  $T_{N1}$  and  $T_{N2}$ . The first one corresponds to the ordering of Cu atoms in the  $\text{CuO}_2$  planes,  $T_{N2}$  reflects the antiferromagnetic ordering of magnetic moments in CuO chains relatively to the moments in the planes.  $T_{N1}$  and  $T_{N2}$  depend strongly on the oxygen content [3]

( $T_{N1}=450$  K for  $x=0,1$  and  $T_{N2}=80$  K, but  $T_{N1}=230$  K and  $T_{N2}=10$  K for  $x=0.35$ ).

We have tried to describe magnetic interactions in high- $T$  superconductors basing on the LMO band structure calculations. Exchange interaction parameters can be defined from the effective Heisenberg hamiltonian:

$$H_{ex} = -1/2 \sum_{ij} J_{ij} \vec{S}_i \cdot \vec{S}_j \quad (1)$$

When the magnetic moments are not too large, as copper magnetic moments in superconducting oxides,  $J_{ij}$  parameters can be defined through the non-local magnetic susceptibility of spin-restricted solution for the crystal [4,5]:

$$J_{ij} = \frac{1}{2} \frac{I_i \cdot I_j}{S_i \cdot S_j} \sum_{LL'} \chi_{LL'}^{ij} \quad (2)$$

where

$$\chi_{LL'}^{ij} = 1/\pi \int_{-\infty}^{\infty} \text{Im } G_{LL'}^{ij}(E) \cdot G_{LL'}^{ji}(E) dE \quad (3)$$

$$G_{LL'}^{ij} = 1/\Omega_{BZ} \int_{BZ} d\vec{k} \sum_n \frac{\psi_{nl}(\vec{k}) \psi_{nl}^*(\vec{k})}{E - E_n(\vec{k})} e^{i\vec{k} \cdot \vec{R}_{ij}} \quad (4)$$

- is the nondiagonal Green function which can be calculated through the energy spectrum  $E_n(\vec{k})$  and  $\psi_{nl}(\vec{k})$  of the LMTO-hamiltonian.  $I_i$  in formulae (2) is interatomic exchange parameters of atom i. Due to the sharp dependence of nondiagonal Green function  $G_{LL}^{ij}(E)$  on the E the integral (3) should be calculated as contour integral in complex plane.

The results of nonlocal magnetic susceptibility calculations and the values of exchange interaction parameters for  $\text{La CuO}$  and  $\text{YBa}_2\text{Cu}_3\text{O}_7$  systems are given in the Table.

Strong anisotropy of exchange interactions in the ab plane and along the c axis in  $\text{La}_2\text{CuO}_4$  is obviously seen. The value of Neel temperature found agrees well with the experimental data available. In the  $\text{YBa}_2\text{Cu}_3\text{O}_7$  system there is strong antiferromagnetic Cu-O-Cu interaction in the  $\text{CuO}_2$  plane, which results in antiferromagnetic ground state of  $\text{YBa}_2\text{Cu}_3\text{O}_6$ .

Superexchange of Cu1-O4-Cu2 type is antiferromagnetic also, in accordance with the experiment. Using the simplest mean field approximation  $T_N = \frac{1}{3} |J| z S(S+1)$ , where z is the number of the nearest magnetic neighbours, it is possible to estimate Neel temperature values. They are  $T_{N1} = 265-314$  K,  $T_{N2} = 61$  K agree well with the experimental data. Large ferromagnetic moment exchange in Cu1-O4-Cu1 chains ( which follows from NMR experiments [6] also ) does not influence antiferromagnetic ordering, as when  $\phi = 1$  all the chains are broken.

In the planes of "1-2-3" system there are quite strong antiferromagnetic Cu-O and O-O interaction which appear due to holes in oxygen subbands. These results are in line with the magnetic model of oxygen holes pairing in high- $T_c$  superconductors suggested in [7].

We have performed also a number of LMTO spin-polarised calculations for  $\text{GdBa}_2\text{Cu}_3\text{O}_6$  and  $\text{GdBa}_2\text{Cu}_3\text{O}_7$  both for ferromagnetic and antiferromagnetic ordering of magnetic moments. For antiferromagnetic state the energy gap of 0,04 eV is formed at the Fermi level. Spin splitting of Gd f-states equals to 5 eV, and magnetic moment of Gd atoms is  $7.3 \mu_B$ . The estimation of stability parameters [8] leads to the conclusion of more stable antiferromagnetic ordering: the difference of  $J_0$  parameters is 230 K ( $\text{O}_6$ ) and 180 K ( $\text{O}_7$ ). The value of s-f integral in antiferromagnetic phase of  $\text{GdBa}_2\text{Cu}_3\text{O}_7$  appears to be about 20 K, and correspond to small changes of  $T_c$  when Gd atoms enter the crystal lattice.

Table. Exchange interaction parameters for  $\text{YBa}_2\text{Cu}_3\text{O}_7$   
and  $\text{La}_2\text{CuO}_4$  (for  $s=1/2$  and  $I_{\text{Cu}}=0,07 \text{ Ry}$ ,  $I_{\text{O}}=0,11 \text{ Ry}$ )

pair	$\hat{R}_{ij}$	$\chi_{dd(p)}^{ij} (\text{mRy}^{-1})$	$J^{ij} (\text{K})$
$\text{YBa}_2\text{Cu}_3\text{O}_7$			
Cu2-O2-Cu2	(0 1 0)	-50,8	-157
Cu2-O3-Cu2	(1 0 0)	-42,8	-132
Cu1-O4-Cu2	(0 0 1)	-9,9	-31
Cu1-O1-Cu1	(0 1 0)	72,5	225
Cu2-O2	(0,5 0 0)	112,3	545
O2-O2	(0,5 0,5 0)	35,5	270
$\text{La}_2\text{CuO}_4$			
Cu-O1-Cu	(1 0 0)	-73,5	227
Cu-O2-Cu	(0 0 1)	-0,45	1,4

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